Discussion. Final atomic parameters are given in Table 1 , bond lengths and angles in Tables 2 and 3, and equations for the least-squares planes in Table 4. The atom numbering is shown in Fig. 1.

Table 4. Equations for the least-squares planes Plane
(A) $\mathrm{C}(1)$ to $\mathrm{C}(10)$ : $0.0164 x+0.1217 y+0.9924 z=1.8598$
(B) $\quad \mathrm{O}(1), \mathrm{C}(11)$ and $\mathrm{C}(12)$ : $0.0625 x+0.0478 y+0.9975 z=1.9606$
(C) $\mathrm{C}(1)$ to $\mathrm{C}(12)$ and $\mathrm{O}(1)$ : $0.0191 x+0.1013 y+0.9947 z=1.7721$
(D) $\quad \mathrm{C}(1)$ to $\mathrm{C}(13)$ and $\mathrm{O}(1)$ :
$0.0368 x+0 \cdot 1091 y+0.9934 z=1.8601$
(E) $\quad \mathrm{C}(1)$ to $\mathrm{C}(10)$ and $\mathrm{C}(13)$ :
$0 \cdot 0258 x+0 \cdot 1300 y+0.9912 z=1.9264$
$\mathrm{C}(13), \mathrm{C}(14)$ and $\mathrm{O}(2)$
$0.2668 x+0.9510 y-0.1591 z=1.7203$
The direction cosines of the normals to various planes
Normal to plane

|  |  | $m$ | $n$ |
| :--- | :---: | :---: | :---: |
| $(A)$ | 0.0164 | 0.1217 | 0.9924 |
| $(B)$ | 0.0625 | 0.0478 | 0.9975 |
| (C) | 0.0191 | 0.1013 | 0.9947 |
| (D) | 0.0368 | 0.1091 | 0.9934 |
| (E) | 0.0258 | 0.1300 | 0.9912 |
| $(F)$ | 0.2668 | 0.9510 | 0.1591 |

The dihedral angle between the planes $(A)$ and $(B)$ is $1^{\circ}$ and between planes $(A)$ and $(F)$ is $86.8^{\circ}$.

The deviations of the atoms in the central aromatic ring are all within experimental error and the ring may be assumed to be planar. The atoms of the ethoxy group are nearly coplanar with the ring atoms, the largest deviation being for $C(12)(0 \cdot 13 \AA) . C(13), C(14)$ and $\mathrm{O}(2)$ of the ethyl methyl ether group are in a plane nearly at right angles to the plane of the aromatic ring.

In the crystal, the molecules are arranged in layers (Fig. 2), the interlayer separation being approximately $3.00 \AA$; they are held together loosely by van der Waals forces, the contacts ranging from 3.04 to $3.98 \AA$. The coefficient of molecular packing (Kitaigorodsky, 1955 ) is 0.72 with a coordination number of six.

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# Hydrogen-Bonded Dimers in Tin(II) Hydrogen Phosphate 

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#### Abstract

SnHPO}_{4}\), monoclinic, $P 2_{1} / c, a=4.608$ (2), $b=13.603$ (4), $c=5.823$ (2) $\AA, \beta=98.76$ (1) ${ }^{\circ}, Z=4$. This neutron diffraction study has shown that the $\mathrm{HPO}_{4}^{2-}$ ions are linked together by two asymmetric hydrogen bonds $[d(\mathrm{O} \cdots \mathrm{O})=2 \cdot 560, d(\mathrm{O}-\mathrm{H})=1.017 \AA, \angle \mathrm{O}-$ $\mathrm{H} \cdots \mathrm{O}=178 \cdot 1^{\circ} \mathrm{]}$ to form dimers. Tin-phosphate-oxygen coordination takes place primarily in layers with the hydrogen bonds occurring between these layers.

Introduction. The structure of $\mathrm{SnHPO}_{4}$ was determined by Berndt \& Lamberg (1971) and further refined by McDonald (1975). Both studies used X-ray diffraction. T. H. Jordan prepared colorless crystals of $\mathrm{SnHPO}_{4}$ by the reaction of $\mathrm{SnCl}_{2}$ and $\mathrm{H}_{3} \mathrm{PO}_{4}$ at $p \mathrm{H} \simeq 2$. The cell dimensions in our study were obtained by least squares

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from the setting angles for 30 reflections measured on a four-circle diffractometer with Mo $K \alpha$ radiation and a graphite monochromator. The crystal used for the collection of neutron data was a diamond-shaped (010) plate with a volume $\sim 3.6 \mathrm{~mm}^{3}$. 1724 reflections within a hemisphere defined by a neutron wavelength of $1.232 \AA$ and a limiting $2 \theta$ of $100^{\circ}$ were measured on a four-circle diffractometer according to the procedure described by Prince (1972).

The heavy-atom parameters as given by Berndt \& Lamberg (1971) were used to phase an $F_{o}$ nuclear density synthesis that revealed the H position. The scattering lengths (in units of $10^{-12} \mathrm{~cm}$ ) used were $\mathrm{Sn}: 0 \cdot 61$, $\mathrm{P}: 0.51, \mathrm{O}: 0.580$ and $\mathrm{H}:-0.374$ (Bacon, 1972). Refinement of atomic parameters was accomplished with the program RFINE 4 (Finger \& Prince, 1975); the sum $\sum w\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2}$, where $w=\sigma^{-2}$, was minimized. The variance of $\left|F_{o}\right|, \sigma^{2}$, was taken as $\sigma_{\text {count }}^{2}+0 \cdot 0004 F_{o}^{2}$.

Equivalent reflections were not averaged and so all reflections in the hemisphere with $\left|F_{o}\right|>2 \sigma$ were used in the refinement. The value of $R\left(=\Sigma| | F_{o}\left|-\left|F_{c}\right|\right| / \Sigma\left|F_{o}\right|\right)$ for these 1473 reflections was 0.079 at convergence. A normal probability plot of the residuals, $\left(\left|F_{o}\right|-\left|F_{c}\right|\right) / \sigma$, against the expected quantiles showed some non-linearity due to outliers. The interquartile range (quantiles -0.667 to +0.667 ) was doubled to give an approximate value of $\Delta_{\max }$ to generate new weights (Nicholson \& Prince, 1977) $w^{\prime}=w\left(1-3 x^{4}+2 x^{6}\right)$ where $x=\left(\left|F_{o}\right|-\left|F_{c}\right|\right) / \sigma \Delta_{\text {max }}$ and $w^{\prime}=0$ if $|x|>1$. Two cycles using $w^{\prime}$ reduced $R$ to 0.073 for 1415 reflections and yielded residuals whose normal probability plot was quite linear. The subsequent interquartile range was multiplizd by 3.5 to yield a value of $\Delta_{\text {max }}$ that changes the weights of $97 \%$ of normally distributed data by less than $9 \%$. Two additional cycles of refinement with the new $\Delta_{\max }=16.5$ rejected only $1 \%$ of the data and gave $R=0.075$ for 1460 reflections. $\dagger$ The normal prob-
$\dagger$ A list of structure factors has been deposited with the British Library Lending Division as Supplementary Publication No. SUP 31914 ( 7 pp .). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 13 White Friars, Chester CH1 1NZ, England.
ability plot of the residuals was essentially linear. Some thermal parameters changed by $2 \sigma$ and the approximate extinction parameter $T r^{*}$ by $4 \sigma$ when weights $w^{\prime}$ were used. The differences between the positional parameters common to this study and a recent study by McDonald (1975) are insignificant. The final atomic parameters are given in Table 1. Selected interatomic distances and angles are given in Table 2.

Discussion. This investigation was undertaken to elucidate the nature of the hydrogen bonding in order to provide a structural basis for the interpretation of some Raman and IR data obtained on $\mathrm{SnHPO}_{4}$. In turn, the spectroscopic characterization of $\mathrm{HPO}_{4}^{2-}$ dimers provides a means of looking for such entities in various situations. For example, the state of $\mathrm{HPO}_{4}^{2-}$ ions in tooth enamel is of interest (Davidson \& Hoekstra, 1975).
Fig. 1 shows the structure as viewed along the $c$ axis. The $\mathrm{HPO}_{4}^{2-}$ dimers lie across centers of symmetry and in sheets parallel to the ( 100 ) plane while the hydrogen bonds are all in planes parallel to ( 010 ). Thus, layers of strong interactions, $\mathrm{Sn} \cdots \mathrm{O}$ bonds, alternate with weak interactions, the hydrogen bonds.

Table 1. Positional and thermal parameters for $\mathrm{SnHPO}_{4}$
The form of the thermal factor is $\exp \left[-2 \pi^{2}\left(\sum_{i j} U_{i j} h_{i} h_{j} a_{i} a_{j}^{*}\right]\right.$. The figures in parentheses are standard deviations as estimated by the last cycle of least-squares refinement. All values are $\times 10^{4}$.

|  | $\boldsymbol{x}$ | $y$ | $z$ | $U_{11}$ | $U_{22}$ | $U_{33}$ | $U_{12}$ | $U_{13}$ | $U_{23}$ |
| :--- | ---: | ---: | ---: | :---: | :---: | :---: | :---: | ---: | ---: |
| Sn | $511(2)$ | $1575(1)$ | $1093(2)$ | $158(6)$ | $183(6)$ | $212(7)$ | $21(3)$ | $-2(4)$ | $-21(4)$ |
| P | $5714(2)$ | $1511(1)$ | $5640(2)$ | $117(6)$ | $150(6)$ | $138(7)$ | $-55(4)$ | $18(5)$ | $-2(4)$ |
| $\mathrm{O}(1)$ | $3027(2)$ | $2129(1)$ | $4566(2)$ | $126(6)$ | $239(7)$ | $185(7)$ | $29(4)$ | $6(5)$ | $-1(4)$ |
| $\mathrm{O}(2)$ | $4807(3)$ | $822(1)$ | $7565(2)$ | $275(7)$ | $207(7)$ | $188(7)$ | $-48(5)$ | $65(5)$ | $11(5)$ |
| $\mathrm{O}(3)$ | $6828(3)$ | $913(1)$ | $3763(2)$ | $242(7)$ | $183(6)$ | $215(7)$ | $-16(4)$ | $87(5)$ | $-50(5)$ |
| $\mathrm{O}(4)$ | $8042(2)$ | $2181(1)$ | $6989(2)$ | $134(6)$ | $197(7)$ | $174(7)$ | $-31(3)$ | $9(4)$ | $-9(4)$ |
| H | $4120(5)$ | $4861(2)$ | $2021(5)$ | $341(12)$ | $310(13)$ | $342(13)$ | $26(9)$ | $83(10)$ | $-34(10)$ |

Table 2. Selected interatomic distances and angles with estimated standard deviations in parentheses

| $i$ | $j$ | $k$ | $d_{i j}(\AA)$ | $d_{j k}(\AA)$ | $d_{i k}(\AA)$ | $\angle i j k\left({ }^{\circ}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{O}(2)$ | H | $\mathrm{O}(3)$ | 1.017 (2) | 1.544 (3) | $2 \cdot 560$ (2) | 178.1 (2) |
| P | $\mathrm{O}(2)$ | H | 1.566 (2) | 1.017 (2) | $2 \cdot 203$ (3) | 115.4 (2) |
| $\mathrm{O}(1)$ | P | $\mathrm{O}(2)$ | 1.548 (2) | 1.566 (2) | $2 \cdot 539$ (2) | 109.3 (1) |
| $\mathrm{O}(1)$ | P | $\mathrm{O}(3)$ | 1.548 (2) | 1.514 (2) | 2.505 (2) | 109.8 (1) |
| $\mathrm{O}(1)$ | P | $\mathrm{O}(4)$ | $1 \cdot 548$ (2) | 1.530 (2) | 2.519 (2) | 109.8 (1) |
| $\mathrm{O}(2)$ | P | $\mathrm{O}(3)$ | 1.566 (2) | 1.514 (2) | 2.532 (2) | $110 \cdot 6$ (1) |
| $\mathrm{O}(2)$ | P | $\mathrm{O}(4)$ | 1.566 (2) | 1.530 (2) | 2.431 (1) | $103 \cdot 4$ (1) |
| O(3) | P | $\mathrm{O}(4)$ | 1.514 (2) | 1.530 (2) | 2.549 (2) | 113.7 (1) |
| $\mathrm{O}(1)$ | Sn | $\mathrm{O}\left(1^{1}\right)^{*}$ | $2 \cdot 300$ (2) | 2.357 (2) | 3.082 (1) | 82.85 (5) |
| $\mathrm{O}(1)$ | Sn | $\mathrm{O}\left(4^{1}\right)$ | 2.300 (2) | $2 \cdot 146$ (2) | 2.713 (2) | $75 \cdot 11$ (6) |
| $\mathrm{O}\left(1^{1}\right)$ | Sn | $\mathrm{O}\left(4^{1}\right)$ | $2 \cdot 357$ (2) | $2 \cdot 146$ (2) | $2 \cdot 874$ (2) | $79 \cdot 15$ (6) |

* Atom $\mathrm{O}\left(n^{i}\right)$ is related to $\mathrm{O}(n)$ by the symmetry operation $x, \frac{1}{2}-y, \frac{1}{2}+z$.



Fig. 1. A stereoscopic pair showing the crystal structure of $\mathrm{SnHPO}_{4}$ as viewed along the $c$ axis. The star marks the origin of the coordinate system.

The observed $\mathrm{O}-\mathrm{H}$ distance is consistent with the curve given by Pimentel \& McClellan (1971) relating $\mathrm{O}-\mathrm{H}$ distance as a function of $\mathrm{O} \cdots \mathrm{O}$ distance. In addition, the $\mathrm{O}-\mathrm{H}$ and $\mathrm{O} \cdots \mathrm{O}$ distances are very close to recent neutron results presented (Fig. 1.4) by Speakman (1972).

A neutron study of $\mathrm{CaHPO}_{4} .2 \mathrm{H}_{2} \mathrm{O}$ (Curry \& Jones, 1971) showed that $\mathrm{HPO}_{4}^{2-}$ ions are linked together by hydrogen bonds forming a chain $\cdots \mathrm{HOPO}_{3}^{2-} \ldots$ $\mathrm{HOPO}_{3}^{2-} \cdots$. The $\mathrm{O} \cdots \mathrm{O}$ distance in this chain is $2.678 \AA$, which is significantly longer than the distance of $2.560 \AA$ for the $\mathrm{HPO}_{4}^{2-}$ dimers in $\mathrm{SnHPO}_{4}$.

Some characteristics of the $\mathrm{HPO}_{4}^{2-}$ ion have been tabulated by Baur (1974). These were obtained from nine published structures in which the dimeric configuration does not occur. Baur gives an average value of $1.590 \AA$ for the $\mathrm{P}-\mathrm{O}(\mathrm{H})$ distance compared to our value of $1.566 \AA$. This is to be expected because the H atom in $\mathrm{SnHPO}_{4}$ is participating in a moderately short, strong hydrogen bond. Baur's data show that $\mathrm{P}-\mathrm{O}(\mathrm{H})$ distances shorten even more ( $1.516 \AA$ ) when the $H$ is involved in a symmetrical hydrogen bond. Thus, the $\mathrm{O} \cdots \mathrm{O}$ distance of $2.560 \AA$, the $\mathrm{O}-\mathrm{H} \cdots \mathrm{O}$ angle of $178.1^{\circ}$, and the $\mathrm{P}-\mathrm{O}(\mathrm{H})$ distance of $1.566 \AA$ in $\mathrm{SnHPO}_{4}$ all indicate that the asymmetric hydrogen bond in dimeric $\left(\mathrm{HPO}_{4}^{2-}\right)_{2}$ is quite strong.

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# Single-Crystal X-ray Study of $\boldsymbol{\beta}$-Uranium Pentafluoride. The Eight Coordination of $\mathbf{U}^{\vee}$ 

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#### Abstract

UF}_{5}\), tetragonal, $I \overline{4} 2 d ; a=11 \cdot 456$ (2), $c=$ $5 \cdot 195$ (1) $\AA ; Z=8$. The structure of $\beta-$ UF $_{5}$ has been solved by a three-dimensional X-ray diffraction study, with counter techniques. Rather than seven-coordinated as originally proposed on the basis of powder data, the $U$ is found to be eight-coordinated, with coordination geometry intermediate between a dodecahedron and a square antiprism.

Introduction. Isolated seven-coordinated fluoride anions of the $d$ transition elements are well established in $\mathrm{NbF}_{7}^{2-}$ and $\mathrm{TaF}_{7}^{2-}$ by X-ray (Hoard, 1939) and neutron (Brown \& Walker, 1966) diffraction. In these $\mathrm{K}_{2} \mathrm{MF}_{7}$ compounds the anion coordination polyhedron is that of a capped trigonal prism [but distorted towards a capped octahedron (Penneman, Ryan \& Rosenzweig, 1973)]. In $\left(\mathrm{NH}_{4}\right)_{3} \mathrm{ZrF}_{7}$ there is evidence for a


pentagonal bipyramid coordination unit (Hurst \& Taylor, 1970), first proposed for this compound and for $\mathrm{K}_{3} \mathrm{UF}_{7}$ by Zachariasen (1954a,b). In both these latter structures there is disorder.

It is interesting to note that there are several actinide elements with radii relatively larger than the $d$ elements mentioned, but whose structures are thought to contain coordination polyhedra of both the above types (see Penneman et al., 1973). However, to date there are no single-crystal studies to substantiate such geometry. Thus, when crystals of $\beta-U F_{5}$ were prepared the opportunity of establishing one of these, i.e. the pentagonal bipyramid geometry for an actinide bound only to fluorine, arose. The dioxopentafluorouranyl case with five equatorial fluorines is well established (Zachariasen, $1954 a$ ). The structure of $\beta-\mathrm{UF}_{5}$ had been deduced by Zachariasen (1949) and has remained a

